

# Electron Beam Stimulated Molecular Motions

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Atomic resolution imaging of single molecules presents an ultimate challenge in the study of molecular nanostructures, including biological macromolecules.<sup>1–3</sup> Transmission electron microscopy is an excellent tool for direct observation of nanostructures at angstrom resolution. However, electron molecular imaging is limited by the undesirable effects of electron beam irradiation.<sup>4–7</sup> During the interaction between incident electrons and sample, part of the incident electrons' energy can be transferred to the sample by exciting bonded electrons or by electron scattering.<sup>8,9</sup> Electron excitation leads to irradiation damage by ionization of individual atoms, generation of phonons, or breaking and forming of chemical bonds in the sample.<sup>10</sup> During electron scattering, energy transferred is dependent on the incident electron energy and the scattering angle.<sup>8</sup> Once the energy transferred exceeds the displacement energy of the atom, knock-on damage takes place. A threshold energy for the knock-on damage was reported to be 100–140 keV for sp<sup>2</sup>-bonded carbon atoms.<sup>8,11</sup> Reducing the electron beam knock-on damage has been the driving force toward the development of lower acceleration voltage, aberration-corrected TEMs.<sup>12,13</sup>

Molecules have a range of chemical bonds, and the intermolecular van der Waals (vdW) bonds are the weakest one. Here, we use C<sub>60</sub>'s encapsulated inside carbon nanotubes (CNTs) (peapods) to study electron beam irradiation effects on vdW bonding. A previous study has shown that the CNT provides an ideal host to investigate the behavior of various molecules in a confined one-dimensional space.<sup>14</sup> For example, several kinds of chemical reactions<sup>15,16</sup> and molecular motions<sup>17–20</sup> have been reported in peapods. This work focuses on the range of molecular motions of the vdW-bonded C<sub>60</sub> molecules that can be stimulated by electron beam irradiation.

**ABSTRACT** Electron microscopy with advances in aberration correction has the power to resolve atoms in single molecules. However, its application is limited by electron irradiation induced molecular motions. A better understanding of damage mechanisms is required to achieve the full potential of electron imaging. Here, we report a direct observation of molecular motions stimulated by an electron beam, which allows us to study the breakdown and formation of molecular bonds using C<sub>60</sub>'s encapsulated inside single-walled carbon nanotubes as a model system. An activation energy of 100 s meV is estimated for the observed molecular motions based on van der Waals interactions. We demonstrate that the molecular confinement can significantly increase the electron energy threshold for breaking the vdW bonds.

**KEYWORDS:** C<sub>60</sub> · carbon nanotube · motion · van der Waals interaction · elastic scattering

We report the observation of a single fullerene's jump in a defective zigzag C<sub>60</sub> molecular chain, back and forth translation of a cluster of C<sub>60</sub>'s in a partially filled CNT, pickup of C<sub>60</sub> molecules, and rotation of a zigzag chain of C<sub>60</sub>'s accompanied by deformation of the host CNT. Using a simple model of vdW potential among C<sub>60</sub> molecules and the host CNT, we estimated the activation energies for some of the molecular motions, ranging from ~0.3 to 0.7 eV. These energies allow us to calculate the probability of electron scattering required to activate the molecular motions and to estimate the electron energy and electron beam current density allowed for electron molecular imaging.

All TEM observations were conducted under 80 kV electron beam irradiation. The electron acceleration voltage was selected to avoid the knock-on damage of the C–C sp<sup>2</sup> bonds.<sup>21,22</sup> Details of the experiment including the sample preparation are described in the Experimental Method section. Modeling of the interaction among C<sub>60</sub>'s and the host CNT was carried out using the Lennard-Jones (LJ)-type vdW potential,<sup>23</sup> which has shown excellent agreement with a number of experimental studies of vdW bonding in carbon materials.<sup>24,25</sup>

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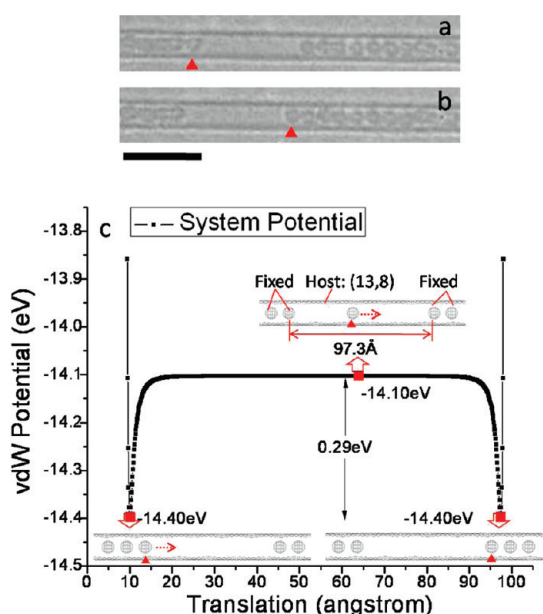


Figure 1. (a, b) Recorded TEM images showing the beginning and ending of a single  $C_{60}$ 's jump in an isolated peapod recorded with an exposure time of 0.5s. The two images are 0.5 s apart. The scale bar is 5 nm. (c) Calculated van der Waals potential energy for the jumping of a single  $C_{60}$  using the models illustrated in the insets.

## RESULTS AND DISCUSSION

**Translation of Single  $C_{60}$ .** Single  $C_{60}$  was found to translate along the tube axis in a partially filled CNT under electron beam irradiation, as shown in Figure 1a,b. The images were captured at 0.5 s apart, which was not fast enough to record the intermediate stages of the molecular jump.<sup>16</sup> Figure 1c plots vdW potential of the peapod system calculated as a function of the moving  $C_{60}$ 's positions. The peapod model used for the vdW potential calculation was constructed using about the same tube diameter as the CNT observed in experiment. The calculations show that the chiral vector of the CNT has very little effect on the potential of a moving  $C_{60}$  inside the CNT.<sup>26,27</sup> Local minima and a local maximum are highlighted in the potential curve with their corresponding structural models. According to Figure 1c, an activation energy around 0.3 eV<sup>19,20,28</sup> is required to break the vdW bonding of a  $C_{60}$  at the end of a  $C_{60}$  molecular chain. After detachment, the  $C_{60}$  can move almost freely inside the CNT.

**Single Molecular Jump.** Figure 2a,b show a marked  $C_{60}$  (by the red triangle) at a defect site of a zigzag molecular chain of  $C_{60}$ 's. During the interval between the two images, the marked  $C_{60}$  jumped from one side to the other side of the chain without introducing obvious effects to the neighboring  $C_{60}$  molecules. Measurement of the inter  $C_{60}$ 's spacing indicates that the  $C_{60}$ 's jump is accompanied by a small translation parallel to the tube axis. Figure 2c plots the contour map of the calculated vdW potential as a function of the  $C_{60}$ 's displacement.  $x$ - and  $y$ -coordinates correspond to the  $C_{60}$ 's displacement parallel and perpendicular to the tube axis, respectively. The lower left and

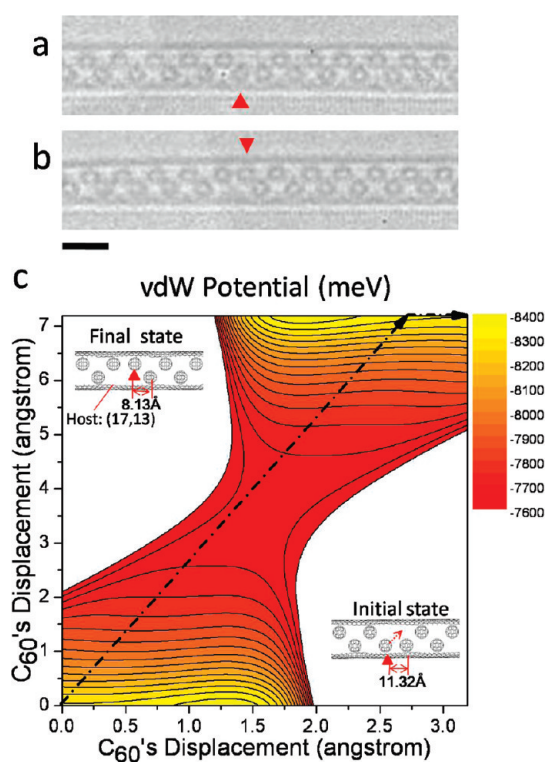


Figure 2. (a, b) Recorded TEM images before and after the jump of a single  $C_{60}$  at a defect site in a zigzag  $C_{60}$  chain. The exposure time is 1 s. The images are 5 s apart. The scale bar is 2 nm. (c) Calculated van der Waals potential energy of the peapod (in meV) as a function of the hopping  $C_{60}$ 's displacement. The  $x$ - and  $y$ -coordinates correspond to the displacement parallel and perpendicular to the tube axis, respectively. Insets are the structural models at the initial and final states used for the calculation.

upper right corners represent the observed initial and final positions of the hopping  $C_{60}$ . In Figure 2c, the calculated energy minima for the initial and final states are deviated from these two corners, which could be due to the thermal energy or excitation of  $C_{60}$  by the electron beam irradiation. One possible path of the  $C_{60}$ 's hopping is marked by the black dashed arrow. Along it, a minimum  $\sim 0.7$  eV in energy is required to drive the motion, which is much higher than the activation energy for the translation of a single  $C_{60}$  as observed in Figure 1.

**Translation and Rotation of a Chain of  $C_{60}$ 's.** A chain of  $C_{60}$ 's performed both translation and rotation inside a partially filled CNT, as shown in Figure 3a–j. Initially, in Figure 3a, a molecular chain of five  $C_{60}$ 's (in the red brackets) was observed at a middle position of the hollow space. Then, the molecular chain moved to the left side, accompanied by an obvious change in its configuration, which appears to be the result of molecular chain's rotation around the tube axis, as shown in Figure 3b. Meanwhile, another  $C_{60}$  (marked by the red triangle in Figure 3b) was attached to the chain. The newly formed chain of six  $C_{60}$ 's continued to change its configuration under the electron beam but without obvious translation during the recording period in Figure 3c,d. In Figure 3e, the molecular chain then moved and got trapped in the middle for about 1 s. After arriving

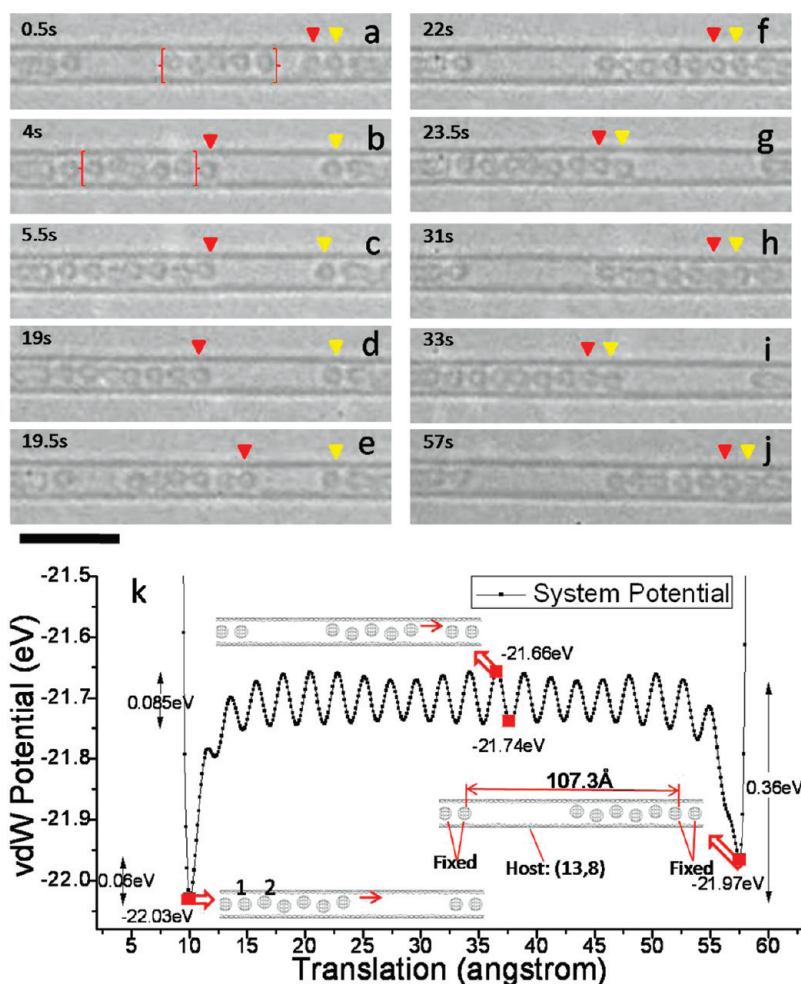


Figure 3. (a–j) Time series of TEM images captured from a video showing back and forth translations of a small chain of  $C_{60}$ 's. The recording lasted for 1 min, and 120 frames in total were recorded with an exposure time of 0.5 s. The scale bar is 5 nm. (k) Calculated van der Waals potential energy for the translation of a chain of five  $C_{60}$ 's inside a CNT. The structural models used for the calculation is shown in insets for three different positions.

at the right side and picking up one more  $C_{60}$  (marked by the yellow triangle) in Figure 3f, a chain of seven  $C_{60}$ 's formed and translated back and forth several times, as captured in Figure 3g–j. Figure 3k plots the calculated vdW potential based on the position of a five- $C_{60}$  molecular chain. To mimic the experimental observation, the  $C_{60}$  chain has a small zigzag arrangement and is allowed to translate along the tube axis. The  $x$ -coordinate represents separation between the two  $C_{60}$  molecules marked as 1 and 2 in the peapod model (lower left in Figure 3k). The estimated activation energy, based on the constructed model, is  $\sim 0.36$  eV (the activation energy is slightly dependent on the  $C_{60}$  orientations, which can introduce a difference of about 0.03 eV). A number of local minima and maxima show up in the middle part of the potential curve with energy barriers of  $\sim 0.085$  eV, which have to be overcome to keep the  $C_{60}$  chain moving continuously inside the single-walled CNT.

**Rotation of a Chain of  $C_{60}$ 's in an Overloaded CNT.** Figure 4a–f show an overloaded peapod with a high line density of  $C_{60}$ 's. In Figure 4a, the  $C_{60}$  chain appears in a compact

and almost linear configuration. Under electron beam irradiation, a part of the  $C_{60}$  molecular chain rotated about the tube axis, with a clear increase in the local projected CNT width, in Figure 4b. The molecular chain was seen in a zigzag arrangement. Separation between neighboring  $C_{60}$ 's in the zigzag direction is comparable to or larger than the CNT's original diameter, which contributes to the deformation seen in the CNT. This observation is followed by another rotation accompanied by a decrease of the CNT projected width in Figure 4c. The swell-and-shrink sequence repeated for a couple of times more, as shown in Figure 4d–f. In Figure 4g, the local width of the CNT around positions A and B, as marked in Figure 4a, was measured at different times. The changes in the observed CNT width can be explained by the elliptical deformation of the host CNT's cross section. As the zigzag  $C_{60}$  chain rotates, the deformation follows, which gives rise to alternate contraction and expansion in projection. The maximum expansion up to 29% was observed in Figure 4d at position A, corresponding to a diameter increase of 0.44 nm.

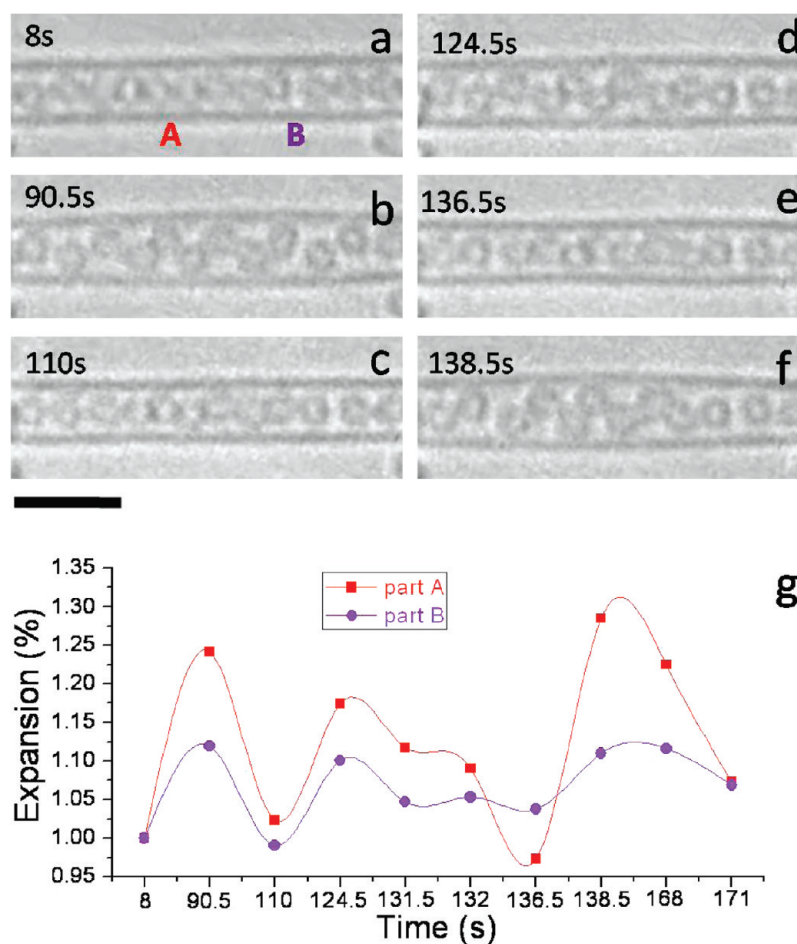


Figure 4. (a–f) Time series of TEM images captured from a video showing rotation of a zigzag C<sub>60</sub> chain leading to expansion and contraction of the host CNT. The scale bar is 5 nm. (g) Measurement of the projected width of the host CNT at positions A and B as marked in (a) at different recording times.

Our work shows that different molecular motions of C<sub>60</sub>'s require different activation energy ranging from  $\sim 0.3$  to  $\sim 0.7$  eV based on the vdW interaction. Somada *et al.*<sup>19</sup> and Ueno *et al.*<sup>29</sup> suggested that the activation energy may come from thermal energy. For this to occur, the temperature has to be as high as 4000 K to activate the single C<sub>60</sub>'s translation in Figure 1. However, the peapod was reported to be an excellent thermal conductor,<sup>30</sup> and an estimation of the sample heating by electron beam irradiation predicted a rise of only less than 1 K above the experimental temperature ( $\sim 300$  K).<sup>31</sup> Thermal gradient driven motion was proposed by Barreiro *et al.*<sup>20</sup> and Zambrano *et al.*<sup>28</sup> In their work, a steady thermal gradient was imposed on the sample to drive the molecular motion through momentum transferred from the phonon excitation.<sup>20</sup> However, in our experiment, considering the lack of temperature rise and relatively uniform illumination during TEM observation,<sup>32</sup> it is unlikely to generate the similar steady thermal gradient in the sample. Moreover, molecular motions driven by the thermal gradient are expected to be directed toward a specific direction, from hot spots to cold regions, which is not consistent

with the back and forth translations nor the rotation of the C<sub>60</sub> molecules.

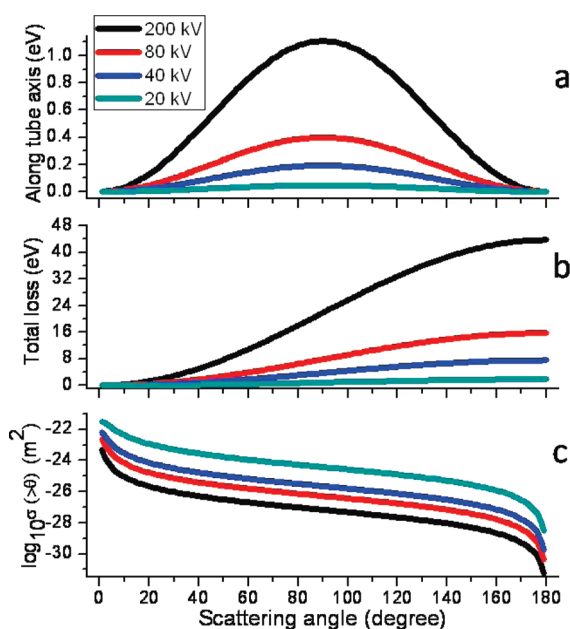
Thus, the activation energy for those observed molecular motions can only come from electron scattering with a significant amount of momentum transfer. In TEM, incident electrons are scattered by atoms in the sample either elastically or inelastically through the Coulomb interaction.<sup>8</sup> Large angle elastic scattering gives rise to large momentum transfer. Because the C<sub>60</sub>'s are confined in a one-dimensional space, only momentum transferred along a specific direction can drive the observed translation. The average momentum transferred from an incident electron to a carbon atom of a C<sub>60</sub> along the tube axis is given by

$$\langle P \rangle = 2 \times \frac{1}{2\pi} \int_0^{\pi/2} m_e \sin \theta \times v_0 \cos \varphi \, d\varphi = \frac{m_e \sin \theta \times v_0}{\pi} \quad (1)$$

corresponding to an amount of energy

$$\langle E_{||} \rangle = \frac{m_e^2 \sin^2 \theta \times v_0^2}{2m_a \pi^2} \quad (2)$$

where  $m_e$ ,  $\theta$ , and  $v_0$  are the mass, scattering angle, and initial velocity of the incident electron,  $m_a$  is the mass of a



**Figure 5.** Energy loss of each incident electron based on the momentum transfer along the tube axis (a) and in an arbitrary direction (total energy loss) (b). (c) Logarithm of the scattering cross section of an incident electron scattered beyond a given scattering angle (x-axis) for different electron acceleration voltages.

carbon atom, and  $\varphi$  is the angle between the tube axis and the plane defined by the momenta of the incident electrons and the carbon atom. Figure 5a plots  $\langle E_{\parallel} \rangle$  as a function of  $\theta$  for different incident electron energies. The maximum  $\langle E_{\parallel} \rangle$  is obtained by scattering the incident electron into  $90^{\circ}$ . In comparison, the total energy transferred,  $\langle E \rangle$ , is

$$\langle E \rangle = \frac{2m_e^2 \sin^2(\theta/2) \cdot v_0^2}{m_a} \quad (3)$$

which has a maximum at  $180^{\circ}$  (Figure 5b).<sup>8</sup>

According to the above equations, for the translation of a  $C_{60}$  molecular cluster as observed in Figure 3a–j, the scattering angle  $\theta$  should be greater than  $72^{\circ}$  in order to transfer the required activation energy ( $\sim 0.36$  eV). The probability of a scattering event is described by its cross section,  $\sigma(\theta)$ , which can be obtained by integrating the differential cross section,<sup>9</sup>

$$\frac{d\sigma(\theta)}{d\Omega} = \frac{\lambda_R^4 Z^2}{64\pi^4 (a_0)^2 (\sin^2(\theta/2) + (\theta_0/2)^2)^2} \quad (4)$$

from  $\theta$  to  $\pi$ . Figure 5c plots the logarithm of  $\sigma(\theta)$  as a function of  $\theta$  under different electron energies. The probability decreases as the scattering angle and the incident electron energy increase. For  $\theta \approx 72^{\circ}$ , the scattering cross section ( $\sigma(\theta) \geq 72^{\circ}$ ) is  $\sim 0.03\%$  of the total scattering cross section ( $2.1 \times 10^{-4} \text{ nm}^2$ ) at 80 kV. With an electron dose of  $3 \times 10^5 \text{ e/s} \cdot \text{nm}^2$ , the chance for a  $C_{60}$  to suffer this kind of scattering can be estimated to be around 1 time per second. During the experiment, the observed frequency of the back and forth translation is 6/min, which is far below the estimated value. Although due to the limited temporal resolution of the microscope, only slow molecular motions are observed.

The maximal energy transfer parallel to the tube axis (eq 2, with  $\theta = 90^{\circ}$ ) is one of  $4\pi^2$  of the maximal total energy that can be transferred (eq 3, with  $\theta = 180^{\circ}$ ). Taking the back and forth translation for example, parallel energy transfer provides a threshold for the incident electron energy at  $72$  kV for  $0.36$  eV loss. In comparison, the total energy transfer has a threshold of  $1.83$  kV for the same energy loss. Thus, for electron molecular imaging, molecular confinement will help the atomic resolution imaging of molecules at medium electron high voltages. Under electron beam irradiation, intermolecular bonding or charge transfer may take place as well to induce the so-called ionization damage. Electron high voltages help reduce the ionization damage.

In summary, we report single  $C_{60}$  jump and translation as well as translation and rotation of a  $C_{60}$  molecular chain. Host CNT expansion up to 29% induced by rotation of a zigzag  $C_{60}$  chain has also been observed. On the basis of van der Waals potential simulation of simple models, we show that these observed molecular motions require  $\sim 0.3$  to  $\sim 0.7$  eV. The lowest activation energy is associated with a single  $C_{60}$  detachment at the end of a  $C_{60}$  molecular chain, while the largest activation energy is associated with the molecular jump at a defect site. This result suggests molecular vacancies can lead to lower activation energy and one-dimensional confinement can significantly increase the electron beam energy threshold for breaking the bonds. For imaging confined  $C_{60}$ 's, the optimum electron energy predicted is  $\sim 60$  kV.

## EXPERIMENTAL METHOD

Single-walled carbon nanotubes used in our experiments were synthesized by the arc discharge method. A mixture of SWCNT and  $C_{60}$  powder was put into a glass tube connected to a vacuum pump. During pumping, the glass tube was heated to  $573$  K for 2 h to get rid of the gas molecules adsorbed on the SWCNT, and then the glass tube was sealed at  $10^{-4}$  Pa and gradually heated and kept at  $723$  K for 48 h. The highly filled

peapods were obtained and then dispersed in *N,N*-dimethyl formamide for ultrasonic treatment. A lacey carbon film on a copper grid obtained from Ted Pella Inc. was used to support the peapods for TEM observation.

To avoid knock-on damage, all TEM observations in this work were performed using a JEOL 2200FS TEM with a field emission gun at a voltage of 80 kV and an electron beam dose of  $3 \times 10^5 \text{ e/s} \cdot \text{nm}^2$  at room temperature.

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**Supporting Information Available:** Movies showing back and forth translations and rotation of C<sub>60</sub> chains. This material is available free of charge via the Internet at <http://pubs.acs.org>.

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